



**Institut für Theoretische Physik**

**Universität Würzburg**

# **Semi-fermionic approach for quantum spin systems**

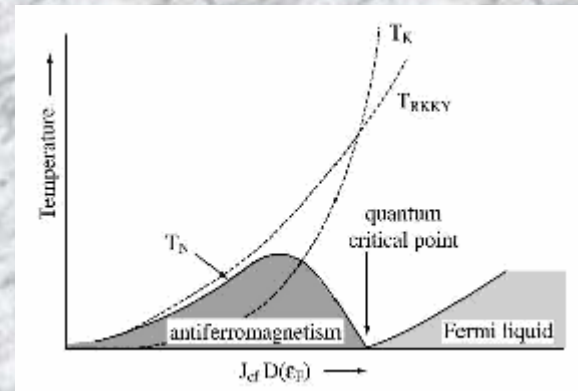
**M.N.Kiselev**



**International Conference on Theoretical Physics  
Paris, UNESCO, 22-27 July 2002**

# Outline

- Motivation
- Model
- Results
- Conclusions



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## References:

- M.Kiselev and R.Oppermann, Schwinger-Keldysh Semionic Approach for Quantum Spin Systems. Phys. Rev. Lett. **85**, 5631 (2000)
- M.Kiselev, H.Feldmann and R.Oppermann, Semi-fermionic representation of SU(N) Hamiltonians. Eur. Phys. J. **B 22**, 53 (2001)
- M.Kiselev, K.Kikoin and R.Oppermann, Ginzburg-Landau functional for nearly Antiferromagnetic perfect and disordered Kondo lattices. Phys. Rev. **B 65**, 184410 (2002)



# Model

$$H = \sum_k \varepsilon(k) c_{k,\sigma}^+ c_{k,\sigma} + J \sum_i \vec{S}_i \vec{s}_i + \sum_{ij} I_{ij} \vec{S}_i \vec{S}_j$$

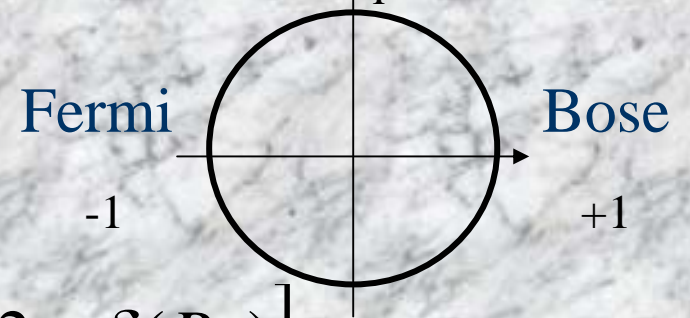
d-electrons
Kondo
RKKY

$$\vec{s} = \frac{1}{2} c_\alpha^+ \vec{\sigma}_{\alpha\beta} c_\beta$$

$$\vec{S} = f_\alpha^+ \vec{\tau}_{\alpha\beta} f_\beta$$

$$S^2 = S(S+1) \Leftrightarrow \sum_\alpha f_\alpha^+ f_\alpha = 1$$

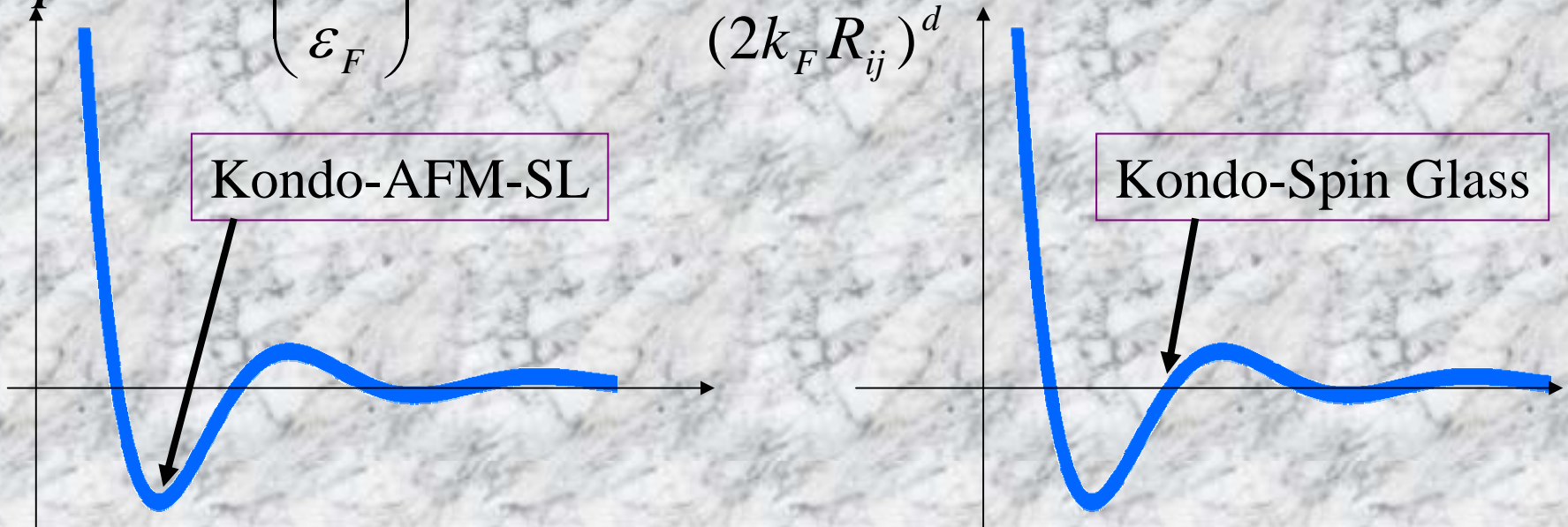
Semi-Fermion



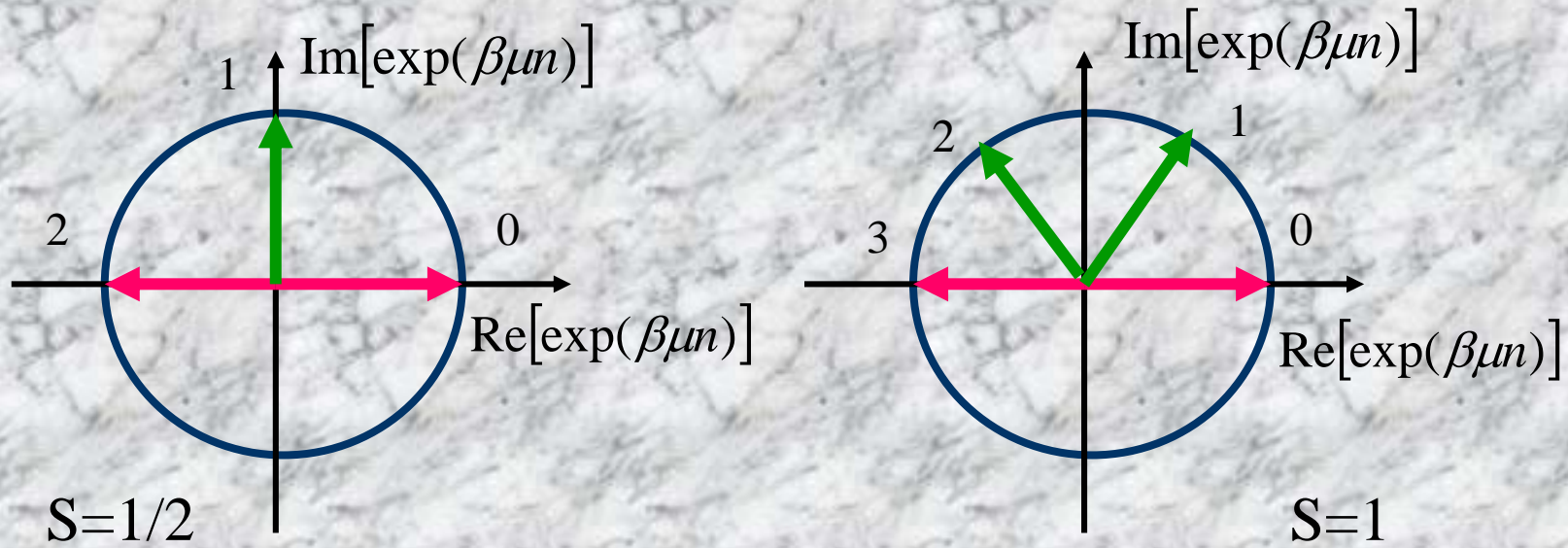
$$I_{ij} = I^{RKKY} = - \left( \frac{J^2}{\varepsilon_F} \right) \frac{\cos[2k_F R_{ij} - \pi(d+1)/2 + \delta(R_{ij})]}{(2k_F R_{ij})^d}$$

Kondo-AFM-SL

Kondo-Spin Glass



# Semi-fermionic representation



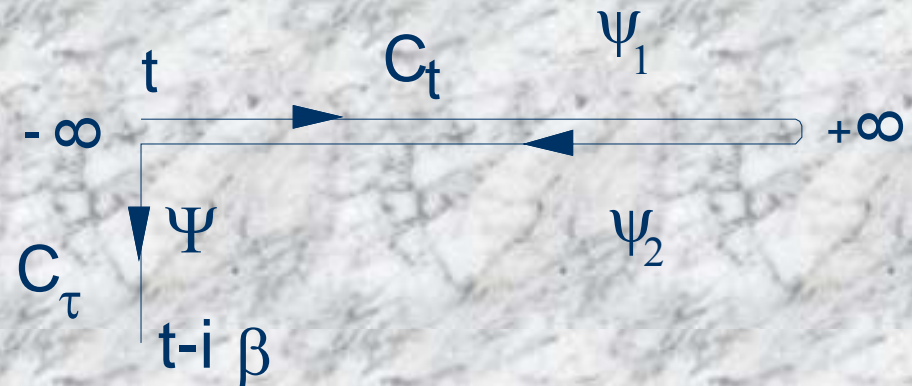
$$\omega = 2\pi T(n + 1/4)$$

$$\omega = 2\pi T(n + 1/3)$$

$$\mu = -i \frac{\pi T}{2S + 1}$$

$$Z_S = \text{Tr}[\exp(-\beta H_S)] = A \text{Tr}[\exp(-\beta H_F + \beta \mu N_F)]$$

$$n^S(\varepsilon) = \frac{1}{\exp(i\pi/(2S+1))\exp(\varepsilon/T) + 1}$$





# RESULTS

Kondo  $\begin{cases} \rightarrow \text{AFM} \\ \rightarrow \text{SL} \\ \rightarrow \text{SG} \end{cases}$

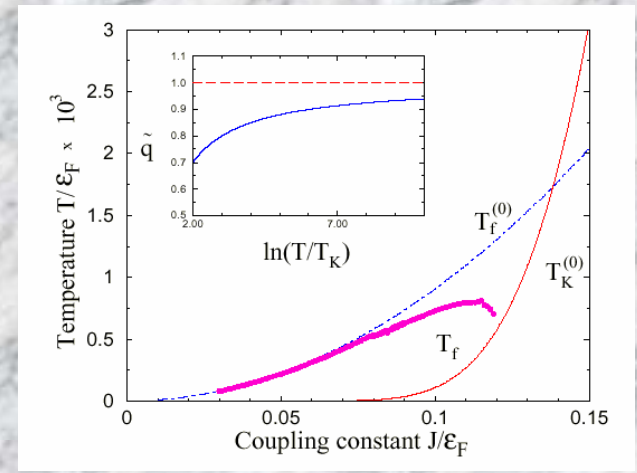
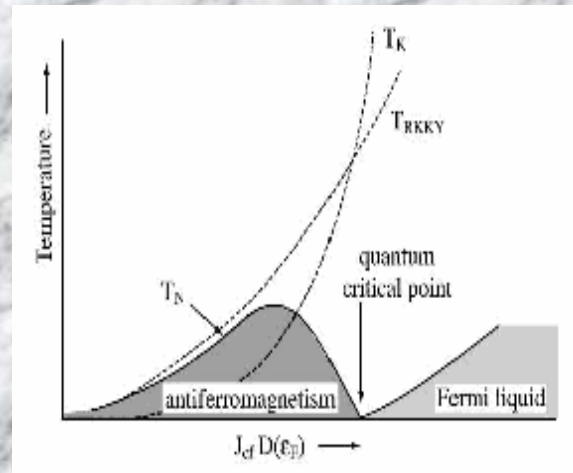
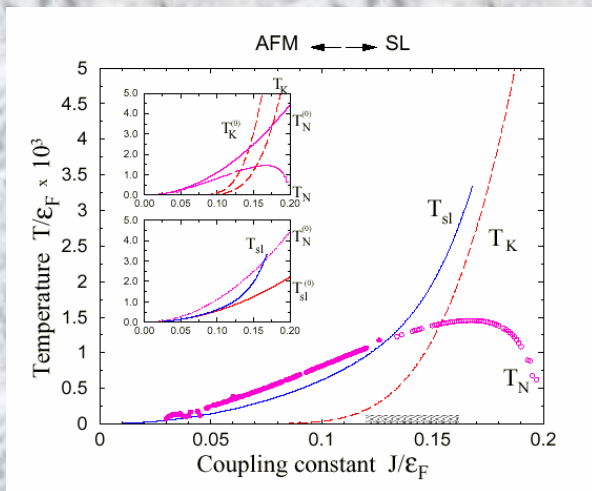
$$A = \sum_{q,n} \left[ \frac{1}{J} - \Pi(N, \Delta, q, \tilde{q}) \right] |\phi(q)|^2 - \text{Tr} \frac{1}{J_Q} N_Q N_{-Q} - \text{Tr} \frac{1}{J_{p-k}} \Delta_p \Delta_k$$

$$N = \tanh \left( \frac{I_Q N}{2T} \right) \left[ 1 - \frac{a_N}{\ln(T/T_K)} \frac{\cosh^2(I_Q N / 2T)}{\cosh^2(I_Q N / T)} \right]$$

$$q = 1 - \frac{c_{sg}}{\ln(T/T_K)}$$

$$\Delta = \sum_q v(q) \left[ \tanh \left( \frac{I_q \Delta}{T} \right) + a_{sl} \frac{I_q \Delta}{T \ln(T/T_K)} \right]$$

$$\tilde{q} = \int_z^G \tanh^2 \left( \frac{I_z \sqrt{q} / T}{1 + 2c_{sg} (I/T)^2 (\tilde{q} - q) / \ln(T/T_K)} \right)$$



# Conclusions

- **Effective action of KL model is written in three-mode representation: Kondo-screening, AFM (SG) and spin-liquid correlations. The problem of local constraint is resolved by means of semi-fermionic representation of the localized spins.**
- **It is shown that the Kondo-screening enhances the tendency of a spin-liquid crossover and partially suppresses AFM and SG orderings.**
- **Competition between local (Kondo) and non-local (magnetic) correlations may result in appearance of Quantum Critical Point.**



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