

Graphene: Relativistic transport in a nearly perfect quantum liquid

Markus Müller

in collaboration with

Lars Fritz (Harvard)

Subir Sachdev (Harvard)

Jörg Schmalian (Iowa)

Sean Hartnoll (Harvard)

Pavel Kovtun (Victoria)

(AdS/CFT)



The Abdus Salam
ICTP Trieste

ICMP 09, Prague 3rd August, 2009

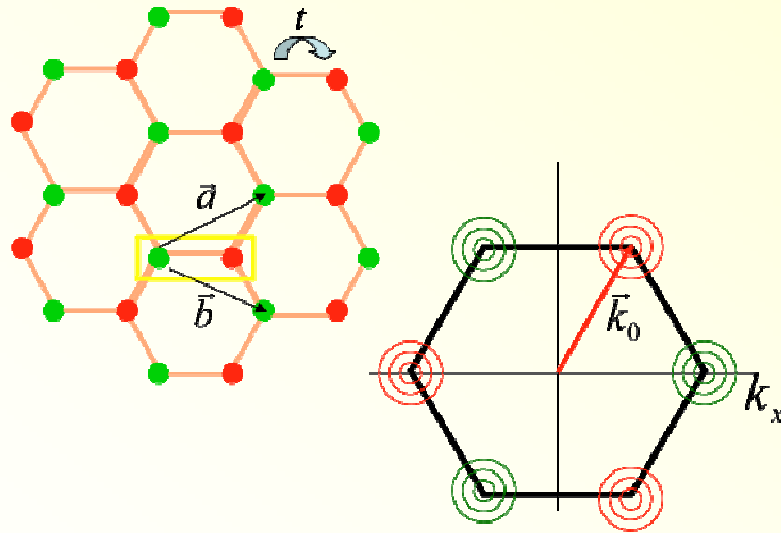
Outline

- Relativistic physics in graphene, quantum critical systems and conformal field theories
- Strong coupling features in collision-dominated transport
- Comparison with strongly coupled fluids (via AdS-CFT)
- Graphene: an almost perfect quantum liquid

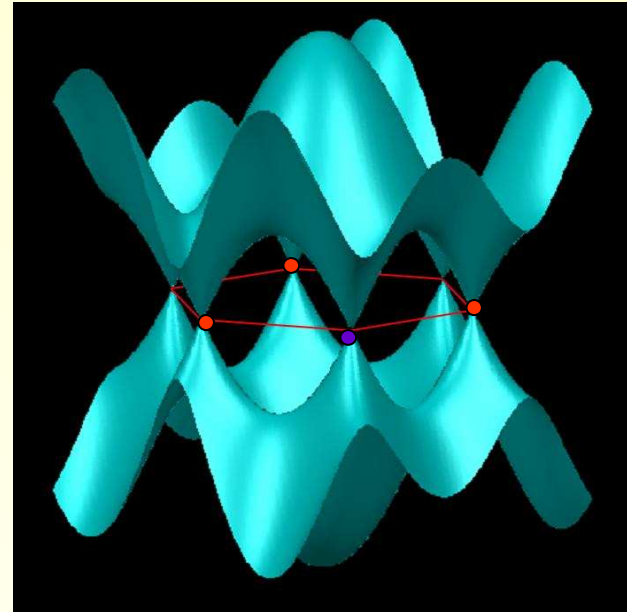
Dirac fermions in graphene

(Semenoff '84, Haldane '88)

Honeycomb lattice of C atoms



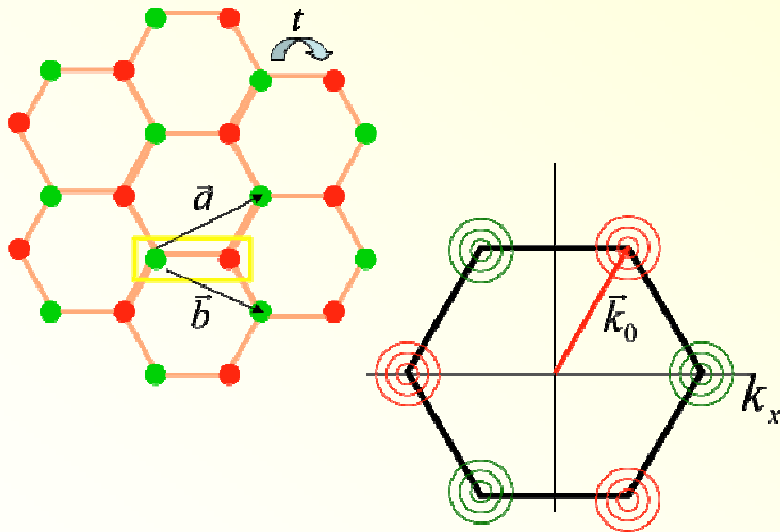
Tight binding dispersion



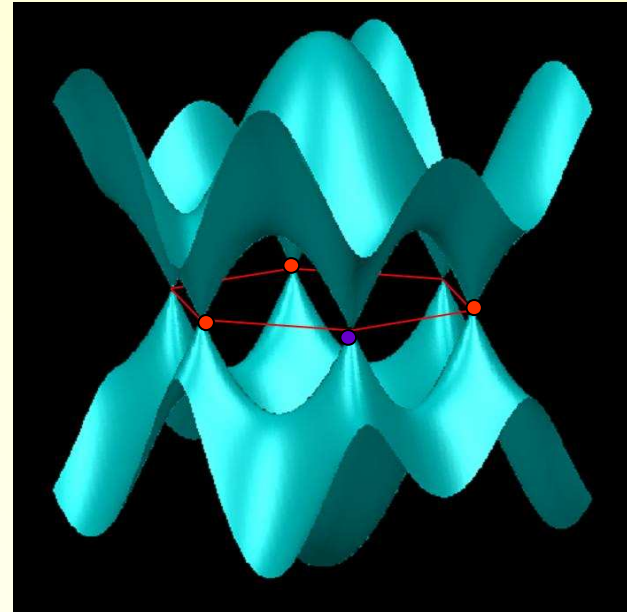
Dirac fermions in graphene

(Semenoff '84, Haldane '88)

Honeycomb lattice of C atoms



Tight binding dispersion



2 massless Dirac cones in the Brillouin zone:
(Sublattice degree of freedom \leftrightarrow pseudospin)

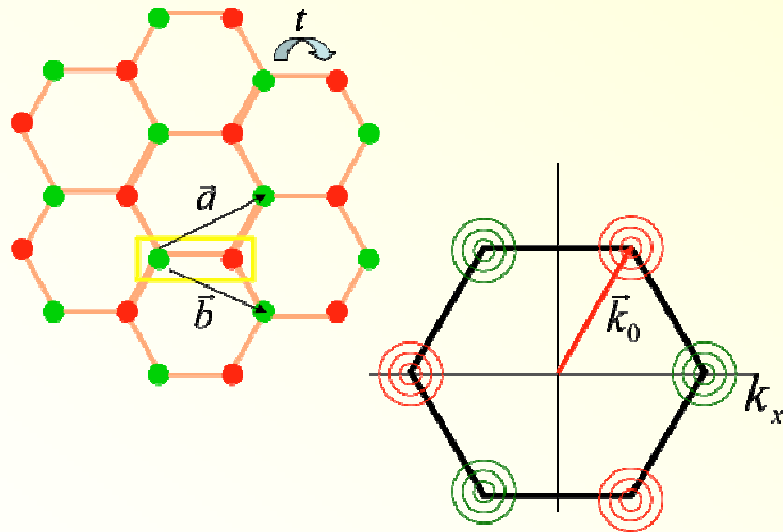
Close to the two Fermi points \mathbf{K}, \mathbf{K}' :

$$H \approx v_F (\vec{\mathbf{p}} - \vec{\mathbf{K}}) \cdot \vec{\sigma}_{\text{sublattice}}$$
$$\rightarrow E_{\mathbf{p}} = v_F |\vec{\mathbf{p}} - \mathbf{K}|$$

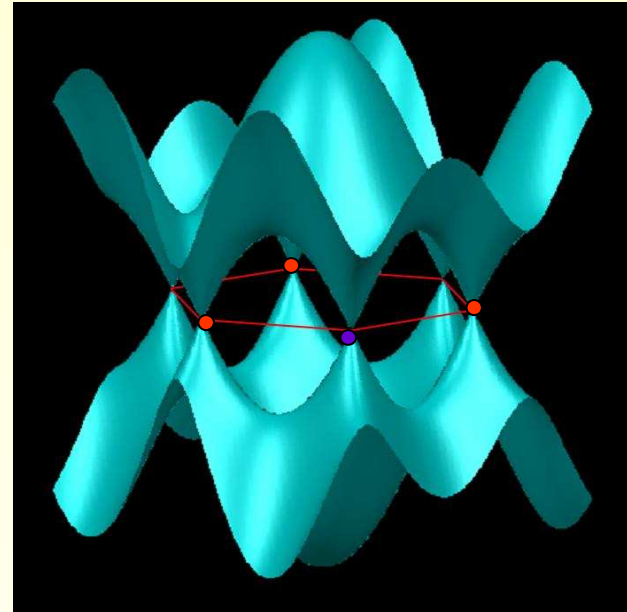
Dirac fermions in graphene

(Semenoff '84, Haldane '88)

Honeycomb lattice of C atoms



Tight binding dispersion



2 massless Dirac cones in the Brillouin zone:
(Sublattice degree of freedom \leftrightarrow pseudospin)

Close to the two Fermi points \mathbf{K}, \mathbf{K}' :

$$H \approx v_F (\vec{\mathbf{p}} - \vec{\mathbf{K}}) \cdot \vec{\sigma}_{\text{sublattice}}$$

$$\rightarrow E_{\mathbf{p}} = v_F |\vec{\mathbf{p}} - \mathbf{K}|$$

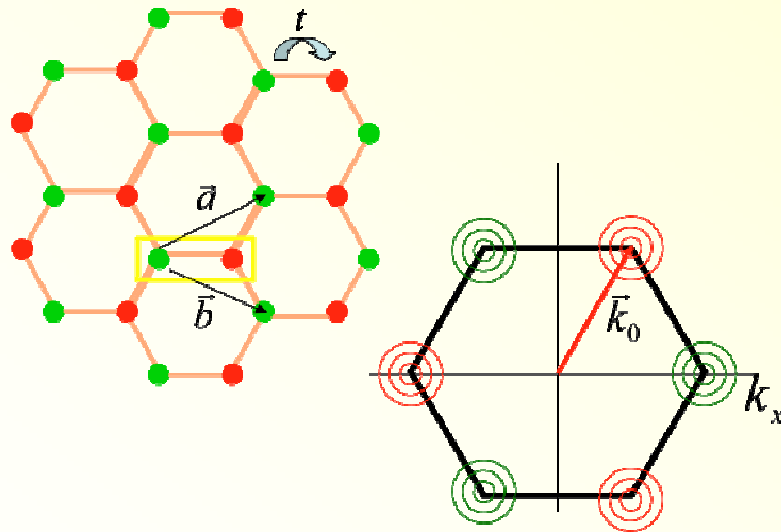
Fermi velocity (speed of light'')

$$v_F \approx 1.1 \cdot 10^6 \text{ m/s} \approx \frac{c}{300}$$

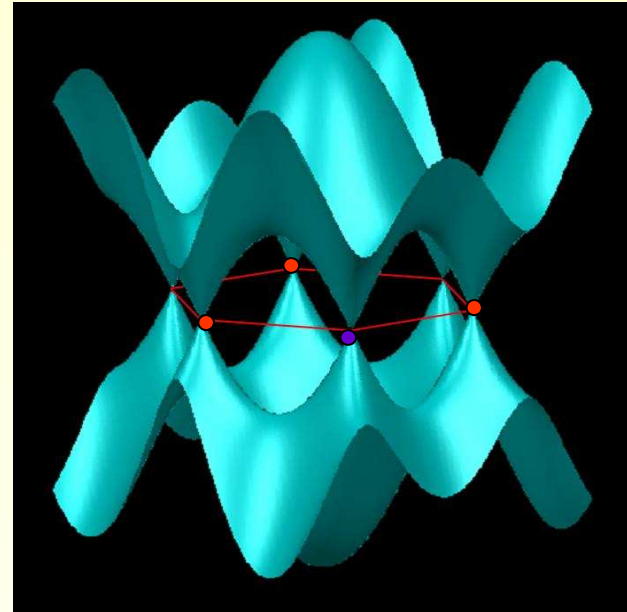
Dirac fermions in graphene

(Semenoff '84, Haldane '88)

Honeycomb lattice of C atoms



Tight binding dispersion



2 massless Dirac cones in the Brillouin zone:
(Sublattice degree of freedom \leftrightarrow pseudospin)

Close to the two Fermi points \mathbf{K}, \mathbf{K}' :

$$H \approx v_F (\vec{\mathbf{p}} - \vec{\mathbf{K}}) \cdot \vec{\sigma}_{\text{sublattice}}$$

$$\rightarrow E_{\mathbf{p}} = v_F |\vec{\mathbf{p}} - \mathbf{K}|$$

Fermi velocity (speed of light'')

$$v_F \approx 1.1 \cdot 10^6 \text{ m/s} \approx \frac{c}{300}$$

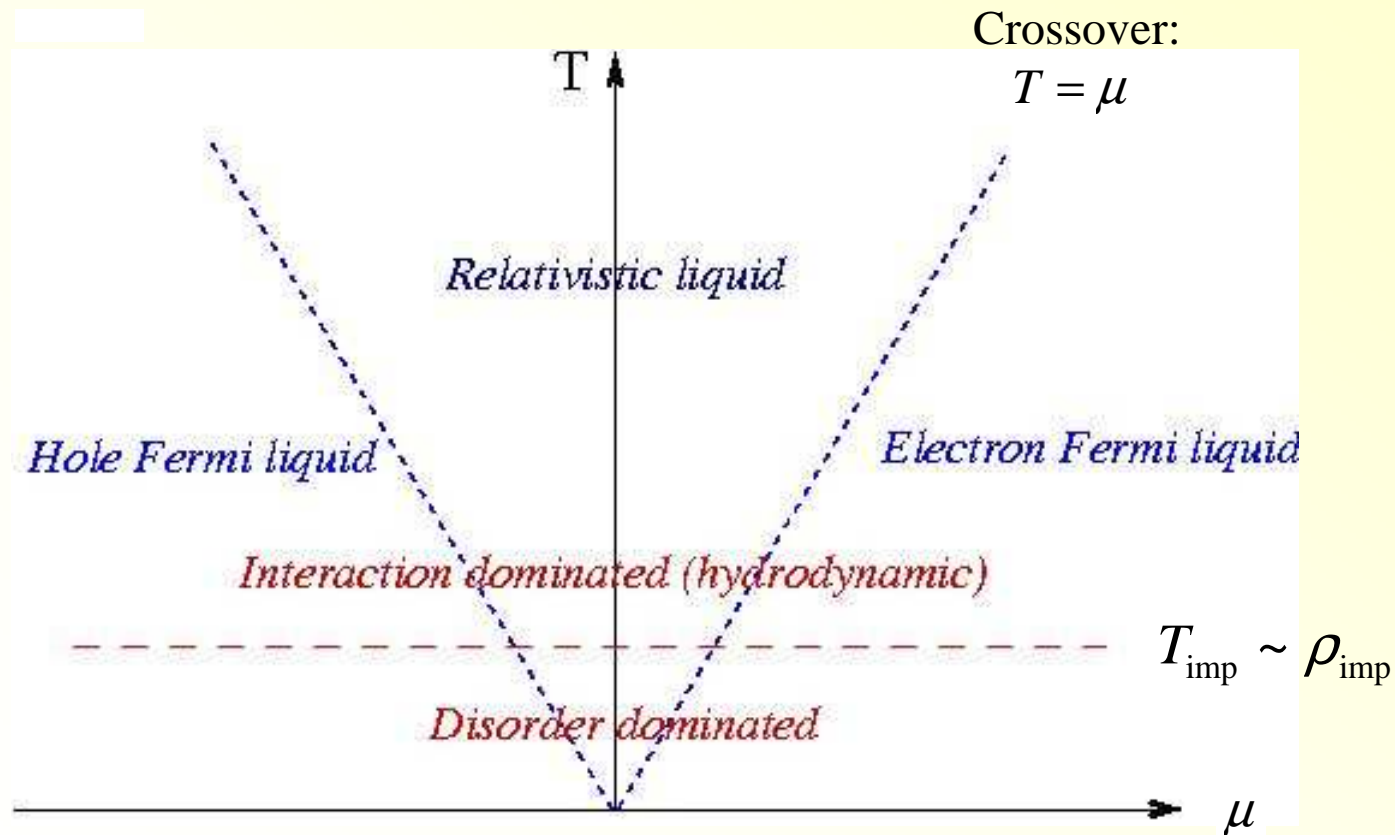
Coulomb interactions: Fine structure constant

$$\alpha \equiv \frac{e^2}{\epsilon \hbar v_F} = O(1)$$

Relativistic fluid at the Dirac point

D. Sheehy, J. Schmalian, Phys. Rev. Lett. 99, 226803 (2007).

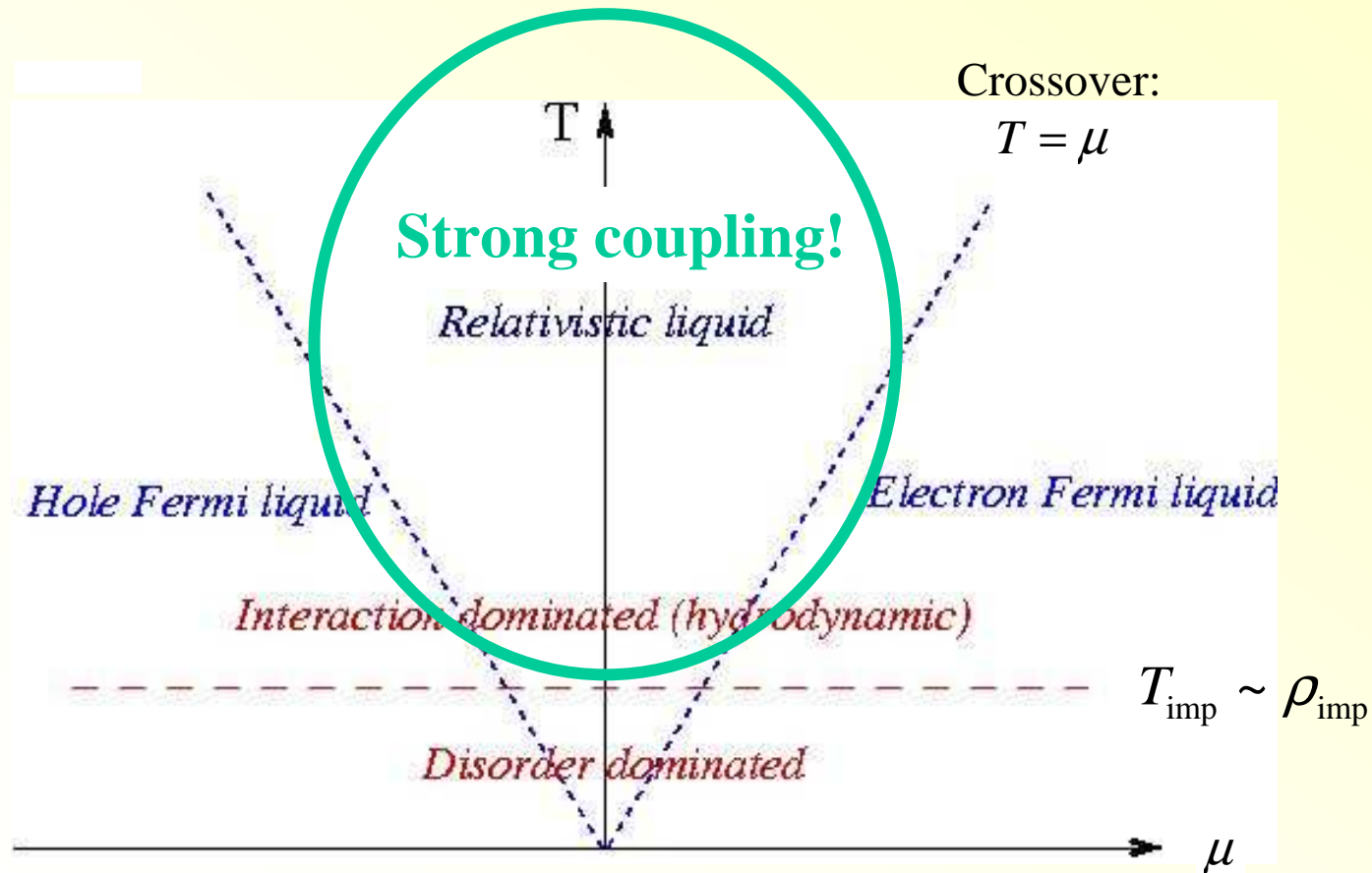
- Relativistic plasma physics of interacting particles and holes!



Relativistic fluid at the Dirac point

D. Sheehy, J. Schmalian, Phys. Rev. Lett. 99, 226803 (2007).

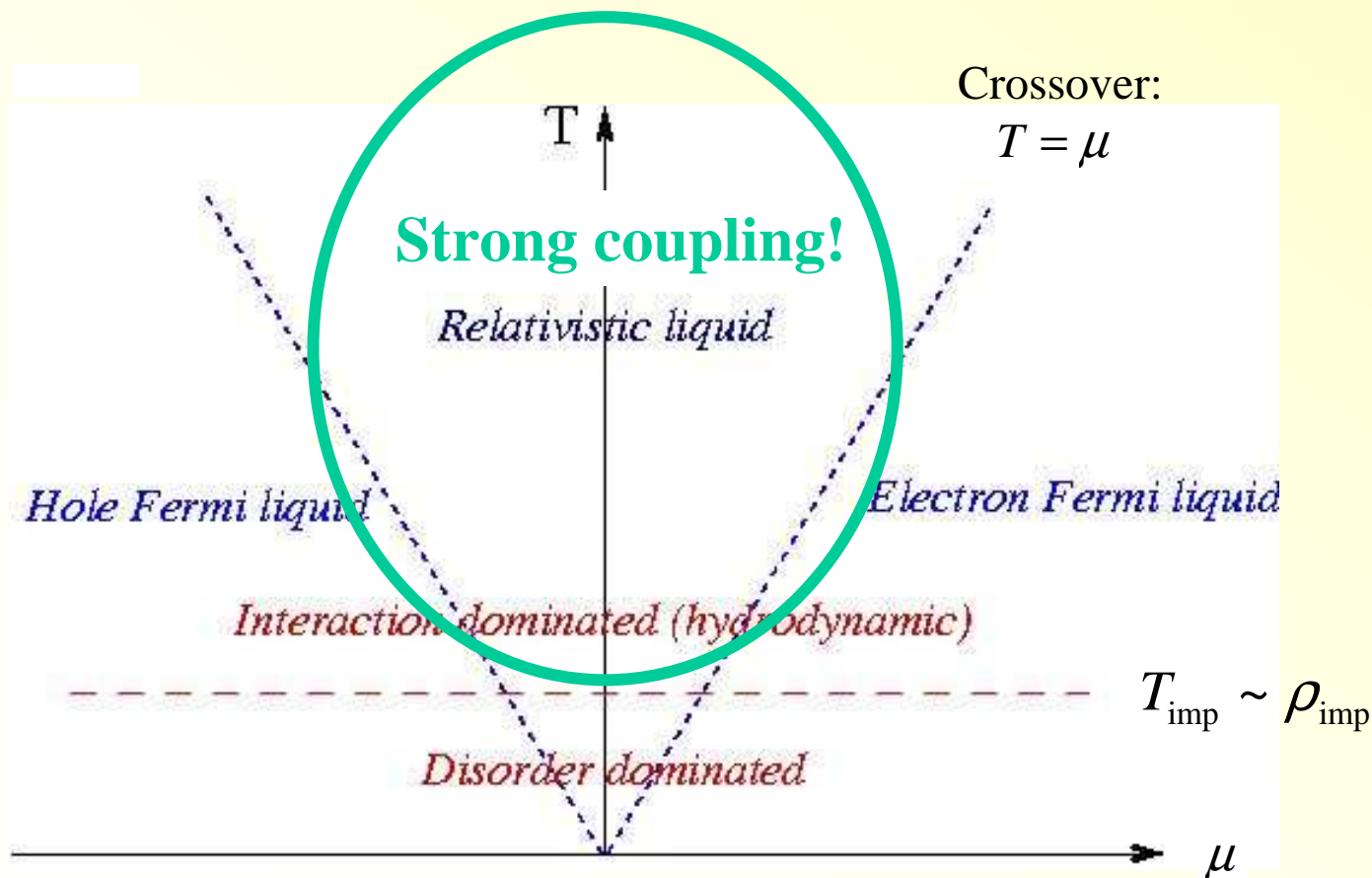
- Relativistic plasma physics of interacting particles and holes!
- Strongly coupled, nearly quantum critical fluid at $\mu = 0$



Relativistic fluid at the Dirac point

D. Sheehy, J. Schmalian, Phys. Rev. Lett. 99, 226803 (2007).

- Relativistic plasma physics of interacting particles and holes!
- Strongly coupled, nearly quantum critical fluid at $\mu = 0$



Very similar as at quantum criticality (with $z=1$, e.g. SIT) and the associated CFT's

Graphene – Fermi liquid?

1. Tight binding kinetic energy
→ massless Dirac quasiparticles

$$H_0 = \sum_{\lambda=\pm} \sum_{a=1}^N \int \frac{d^2 k}{(2\pi)^2} \lambda v_F k \gamma_{\lambda a}^\dagger(\mathbf{k}) \gamma_{\lambda a}(\mathbf{k})$$

2. Coulomb interactions:

Unexpectedly strong!

→ nearly quantum critical!

$$V(\mathbf{q}) = \frac{2\pi e^2}{\epsilon|\mathbf{q}|}$$

$$H_1 = \frac{1}{2} \int \frac{d^2 k_1}{(2\pi)^2} \frac{d^2 k_2}{(2\pi)^2} \frac{d^2 q}{(2\pi)^2} \Psi_a^\dagger(\mathbf{k}_2 - \mathbf{q}) \Psi_a(\mathbf{k}_2) V(\mathbf{q}) \Psi_b^\dagger(\mathbf{k}_1 + \mathbf{q}) \Psi_b(\mathbf{k}_1)$$

Graphene – Fermi liquid?

1. Tight binding kinetic energy
→ massless Dirac quasiparticles

$$H_0 = \sum_{\lambda=\pm} \sum_{a=1}^N \int \frac{d^2 k}{(2\pi)^2} \lambda v_F k \gamma_{\lambda a}^\dagger(\mathbf{k}) \gamma_{\lambda a}(\mathbf{k})$$

2. Coulomb interactions:

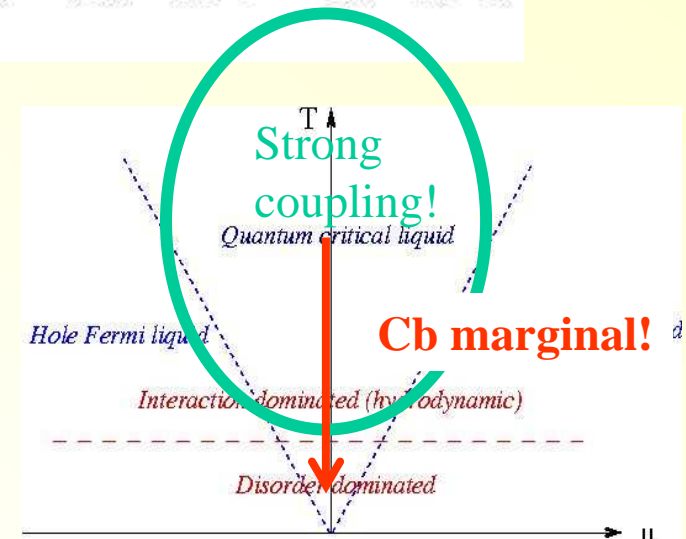
Unexpectedly strong!

→ nearly quantum critical!

$$V(\mathbf{q}) = \frac{2\pi e^2}{\epsilon|\mathbf{q}|}$$

$$H_1 = \frac{1}{2} \int \frac{d^2 k_1}{(2\pi)^2} \frac{d^2 k_2}{(2\pi)^2} \frac{d^2 q}{(2\pi)^2} \Psi_a^\dagger(\mathbf{k}_2 - \mathbf{q}) \Psi_a(\mathbf{k}_2) V(\mathbf{q}) \Psi_b^\dagger(\mathbf{k}_1 + \mathbf{q}) \Psi_b(\mathbf{k}_1)$$

Coulomb **only marginally** irrelevant for $\mu = 0$!



Graphene – Fermi liquid?

1. Tight binding kinetic energy
→ massless Dirac quasiparticles

$$H_0 = \sum_{\lambda=\pm} \sum_{a=1}^N \int \frac{d^2 k}{(2\pi)^2} \lambda v_F k \gamma_{\lambda a}^\dagger(\mathbf{k}) \gamma_{\lambda a}(\mathbf{k})$$

2. Coulomb interactions:

Unexpectedly strong!

→ nearly quantum critical!

$$V(\mathbf{q}) = \frac{2\pi e^2}{\epsilon |\mathbf{q}|}$$

$$H_1 = \frac{1}{2} \int \frac{d^2 k_1}{(2\pi)^2} \frac{d^2 k_2}{(2\pi)^2} \frac{d^2 q}{(2\pi)^2} \Psi_a^\dagger(\mathbf{k}_2 - \mathbf{q}) \Psi_a(\mathbf{k}_2) V(\mathbf{q}) \Psi_b^\dagger(\mathbf{k}_1 + \mathbf{q}) \Psi_b(\mathbf{k}_1)$$

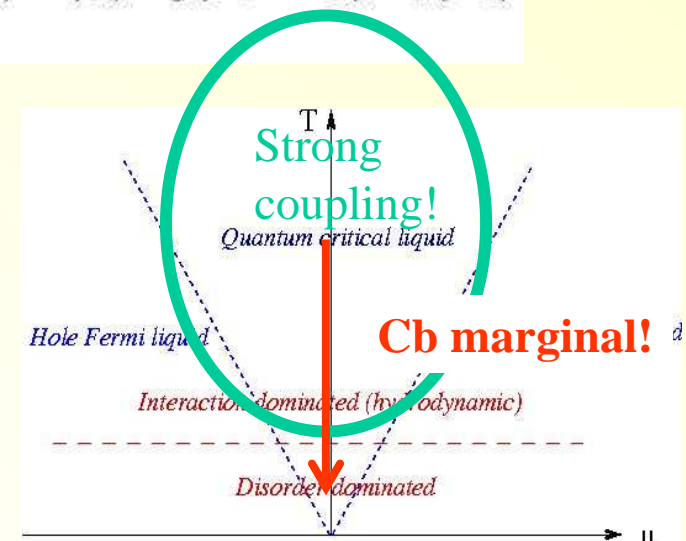
Coulomb **only marginally** irrelevant for $\mu = 0$!

RG:
($\mu = 0$)

$$\frac{d\alpha}{dl} = -\frac{\alpha^2}{4} + \mathcal{O}(\alpha^3)$$

$$\alpha \equiv \frac{e^2}{\epsilon \hbar v_F} = \mathcal{O}(1)$$

$$\alpha(T) = \frac{\alpha^0}{1 + (\alpha^0/4) \ln(\Lambda/T)} \stackrel{T \rightarrow 0}{\sim} \frac{4}{\ln(\Lambda/T)}$$



Graphene – Fermi liquid?

1. Tight binding kinetic energy
→ massless Dirac quasiparticles

$$H_0 = \sum_{\lambda=\pm} \sum_{a=1}^N \int \frac{d^2 k}{(2\pi)^2} \lambda v_F k \gamma_{\lambda a}^\dagger(\mathbf{k}) \gamma_{\lambda a}(\mathbf{k})$$

2. Coulomb interactions:

Unexpectedly strong!

→ nearly quantum critical!

$$V(\mathbf{q}) = \frac{2\pi e^2}{\epsilon |\mathbf{q}|}$$

$$H_1 = \frac{1}{2} \int \frac{d^2 k_1}{(2\pi)^2} \frac{d^2 k_2}{(2\pi)^2} \frac{d^2 q}{(2\pi)^2} \Psi_a^\dagger(\mathbf{k}_2 - \mathbf{q}) \Psi_a(\mathbf{k}_2) V(\mathbf{q}) \Psi_b^\dagger(\mathbf{k}_1 + \mathbf{q}) \Psi_b(\mathbf{k}_1)$$

Coulomb **only marginally** irrelevant for $\mu = 0$!

RG:

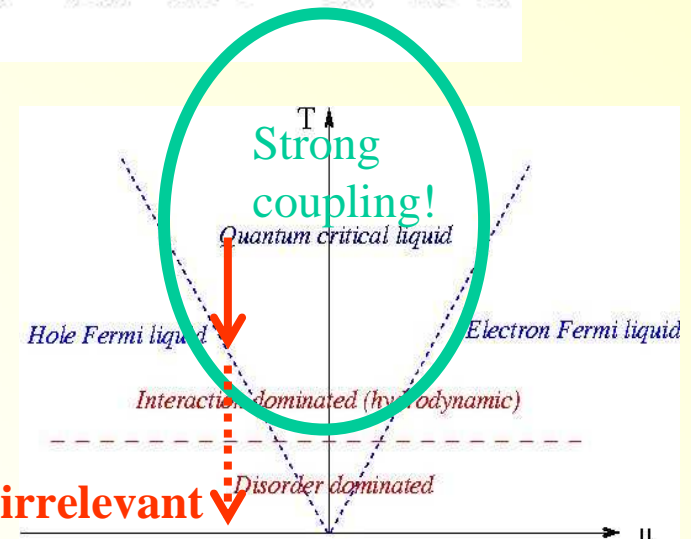
$$\frac{d\alpha}{dl} = -\frac{\alpha^2}{4} + \mathcal{O}(\alpha^3)$$

$$\alpha(T) = \frac{\alpha^0}{1 + (\alpha^0/4) \ln(\Lambda/T)} \stackrel{T \rightarrow 0}{\sim} \frac{4}{\ln(\Lambda/T)}$$

$\alpha \equiv \frac{e^2}{\epsilon \hbar v_F} = \mathcal{O}(1)$

$(\mu = 0)$

$(\mu > 0)$ $T < \mu$: Screening kicks in, short ranged Cb irrelevant



Strong coupling in undoped graphene

MM, L. Fritz, and S. Sachdev, PRB '08.

Inelastic scattering rate
(Electron-electron interactions)

$\mu > T$: standard 2d
Fermi liquid

$$\tau_{ee}^{-1} \sim \alpha^2 \frac{k_B T^2}{\hbar \mu}$$

Strong coupling in undoped graphene

MM, L. Fritz, and S. Sachdev, PRB '08.

Inelastic scattering rate
(Electron-electron interactions)

$\mu > T$: standard 2d
Fermi liquid

$$\tau_{ee}^{-1} \sim \alpha^2 \frac{k_B T^2}{\hbar \mu}$$

Relaxation rate $\sim T$,
like in quantum critical systems!
Fastest possible rate!

$\mu < T$: strongly
coupled relativistic
liquid

$$\tau_{ee}^{-1} \sim \alpha^2 \frac{k_B T}{\hbar}$$

Strong coupling in undoped graphene

MM, L. Fritz, and S. Sachdev, PRB '08.

Inelastic scattering rate
(Electron-electron interactions)

$\mu > T$: standard 2d
Fermi liquid

$$\tau_{ee}^{-1} \sim \alpha^2 \frac{k_B T^2}{\hbar \mu}$$

Relaxation rate $\sim T$,
like in quantum critical systems!
Fastest possible rate!

$\mu < T$: strongly
coupled relativistic
liquid

$$\tau_{ee}^{-1} \sim \alpha^2 \frac{k_B T}{\hbar}$$

“Heisenberg uncertainty principle for well-defined quasiparticles”

$$E_{qp} (\sim k_B T) \geq \Delta E_{\text{int}} = \hbar \tau_{ee}^{-1} \sim \alpha^2 k_B T$$

Strong coupling in undoped graphene

MM, L. Fritz, and S. Sachdev, PRB '08.

Inelastic scattering rate
(Electron-electron interactions)

$\mu > T$: standard 2d
Fermi liquid

$$\tau_{ee}^{-1} \sim \alpha^2 \frac{k_B T^2}{\hbar \mu}$$

Relaxation rate $\sim T$,
like in quantum critical systems!
Fastest possible rate!

$\mu < T$: strongly
coupled relativistic
liquid

$$\tau_{ee}^{-1} \sim \alpha^2 \frac{k_B T}{\hbar}$$

“Heisenberg uncertainty principle for well-defined quasiparticles”

$$E_{qp} (\sim k_B T) \geq \Delta E_{\text{int}} = \hbar \tau_{ee}^{-1} \sim \alpha^2 k_B T$$

As long as $\alpha(T) \sim 1$, energy uncertainty is saturated, scattering is maximal
→ Nearly universal strong coupling features in transport,
similarly as at the 2d superfluid-insulator transition [*Damle, Sachdev (1996, 1997)*]

Consequences for transport

1. Collisionlimited conductivity σ in clean undoped graphene
2. Graphene - a perfect quantum liquid: very small viscosity η !

Consequences for transport

1. Collisionlimited conductivity σ in clean undoped graphene
2. Graphene - a perfect quantum liquid: very small viscosity η !
3. Emergent relativistic invariance at low frequencies!

Despite the breaking of relativistic invariance by

- finite T ,
- finite μ ,
- instantaneous $1/r$ Coulomb interactions

Consequences for transport

1. Collisionlimited conductivity σ in clean undoped graphene
2. Graphene - a perfect quantum liquid: very small viscosity η !
3. Emergent relativistic invariance at low frequencies!

Despite the breaking of relativistic invariance by

- finite T ,
- finite μ ,
- instantaneous $1/r$ Coulomb interactions

Collision-dominated transport \rightarrow relativistic hydrodynamics:

- a) Response fully determined by covariance, thermodynamics, and σ , η
- b) Collective cyclotron resonance in small magnetic field (low frequency)

Hydrodynamic regime:
(collision-dominated)

$$\tau_{ee}^{-1} \gg \tau_{imp}^{-1}, \omega_c^{typ}, \omega_{AC}$$

Collision-limited conductivity

Damle, Sachdev, (1996).

Fritz et al. (2008), Kashuba (2008)

Finite conductivity in a pure system at particle-hole symmetry ($\rho = 0$)!

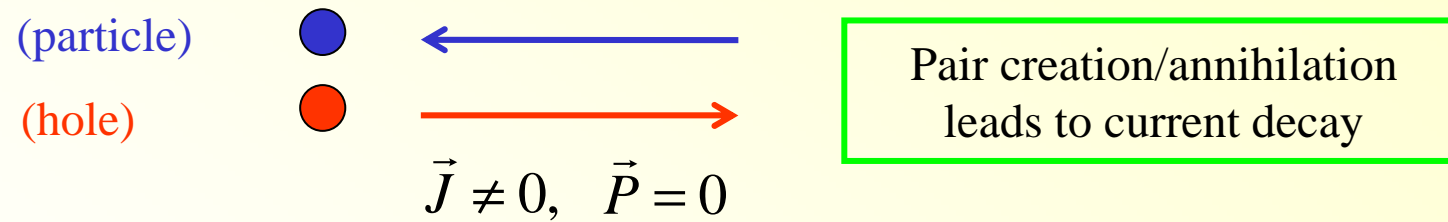
Collision-limited conductivity

Damle, Sachdev, (1996).

Fritz et al. (2008), Kashuba (2008)

Finite conductivity in a pure system at particle-hole symmetry ($\rho = 0$)!

- Key: Charge current without momentum



- Finite collision-limited conductivity! $\sigma(\mu = 0) < \infty$; $\sigma(\mu \neq 0) = \infty$

and

- Infinite thermal conductivity! $\kappa(\mu = 0) = \infty$; $\kappa(\mu \neq 0) < \infty$

(true also in pure semiconductors)

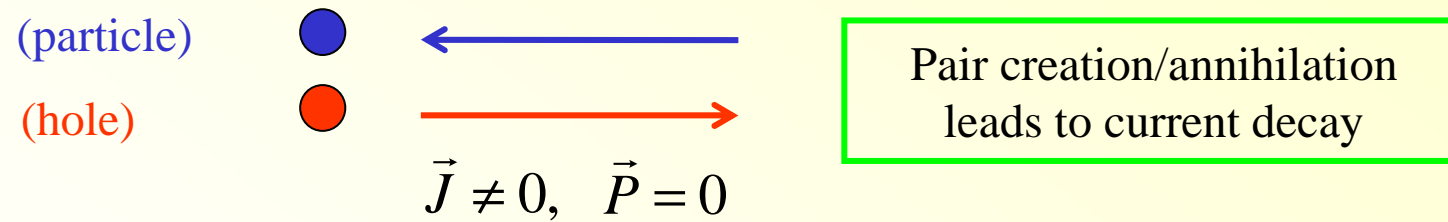
Collision-limited conductivity

Damle, Sachdev, (1996).

Fritz et al. (2008), Kashuba (2008)

Finite conductivity in a pure system at particle-hole symmetry ($\rho = 0$)!

- Key: Charge current without momentum



- **Finite collision-limited conductivity!** $\sigma(\mu = 0) < \infty$; $\sigma(\mu \neq 0) = \infty$
- Only marginal irrelevance of Coulomb:
Maximal possible relaxation rate,
set only by temperature

$$\tau_{ee}^{-1} \approx \alpha^2 \frac{k_B T}{\hbar}$$

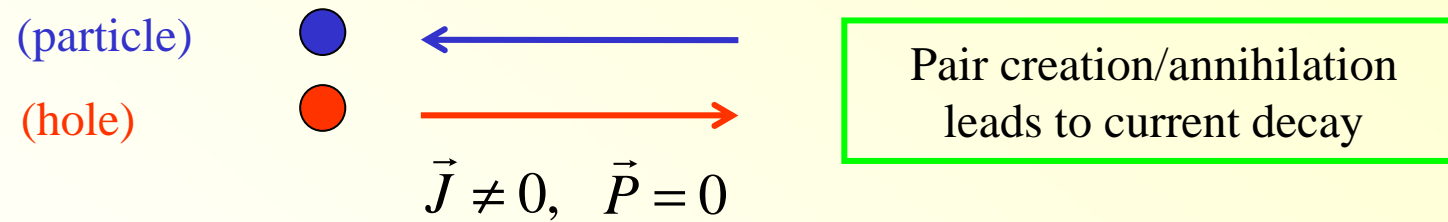
Collision-limited conductivity

Damle, Sachdev, (1996).

Fritz et al. (2008), Kashuba (2008)

Finite conductivity in a pure system at particle-hole symmetry ($\rho = 0$)!

- Key: Charge current without momentum



- **Finite collision-limited conductivity!** $\sigma(\mu = 0) < \infty$; $\sigma(\mu \neq 0) = \infty$
- Only marginal irrelevance of Coulomb:
Maximal possible relaxation rate,
set only by temperature

$$\tau_{ee}^{-1} \approx \alpha^2 \frac{k_B T}{\hbar}$$

→ Nearly universal conductivity at strong coupling

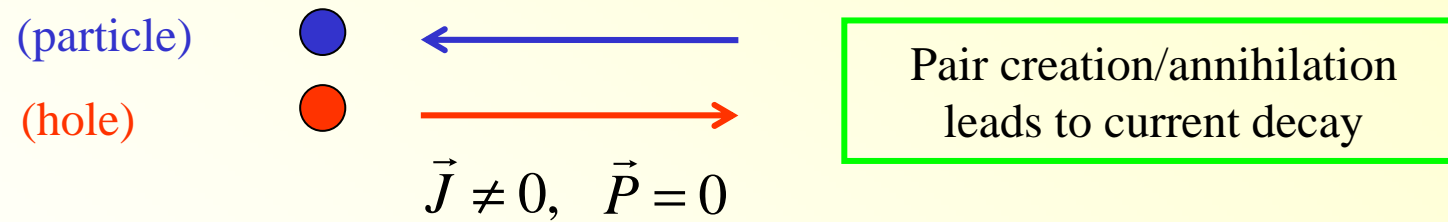
Collision-limited conductivity

Damle, Sachdev, (1996).

Fritz et al. (2008), Kashuba (2008)

Finite conductivity in a pure system at particle-hole symmetry ($\rho = 0$)!

- Key: Charge current without momentum



- **Finite collision-limited conductivity!** $\sigma(\mu = 0) < \infty$; $\sigma(\mu \neq 0) = \infty$
- Only marginal irrelevance of Coulomb:
Maximal possible relaxation rate,
set only by temperature

$$\tau_{ee}^{-1} \approx \alpha^2 \frac{k_B T}{\hbar}$$

→ Nearly universal conductivity at strong coupling

$$\sigma_{Drude} = \frac{e}{m} \rho \tau \rightarrow \sigma(\mu = 0) \sim \frac{e}{k_B T / v^2} \left(e \frac{(k_B T)^2}{(\hbar v)^2} \right) \frac{\hbar}{\alpha^2 k_B T} \sim \frac{1}{\alpha^2} \frac{e^2}{h}$$

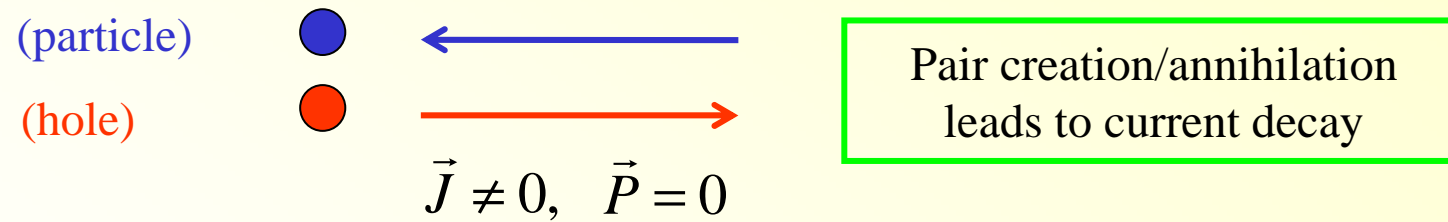
Collision-limited conductivity

Damle, Sachdev, (1996).

Fritz et al. (2008), Kashuba (2008)

Finite conductivity in a pure system at particle-hole symmetry ($\rho = 0$)!

- Key: Charge current without momentum



- **Finite collision-limited conductivity!** $\sigma(\mu = 0) < \infty$; $\sigma(\mu \neq 0) = \infty$
- Only marginal irrelevance of Coulomb:
Maximal possible relaxation rate,
set only by temperature

$$\tau_{ee}^{-1} \approx \alpha^2 \frac{k_B T}{\hbar}$$

→ Nearly universal conductivity at strong coupling

$$\sigma_{Drude} = \frac{e}{m} \rho \tau \rightarrow \sigma(\mu = 0) \sim \frac{e}{k_B T / v^2} \left(e \frac{(k_B T)^2}{(\hbar v)^2} \right) \frac{\hbar}{\alpha^2 k_B T} \sim \frac{1}{\alpha^2} \frac{e^2}{h}$$

Marginal irrelevance of Coulomb:

$$\alpha \approx \frac{4}{\log(\Lambda/T)} < 1$$

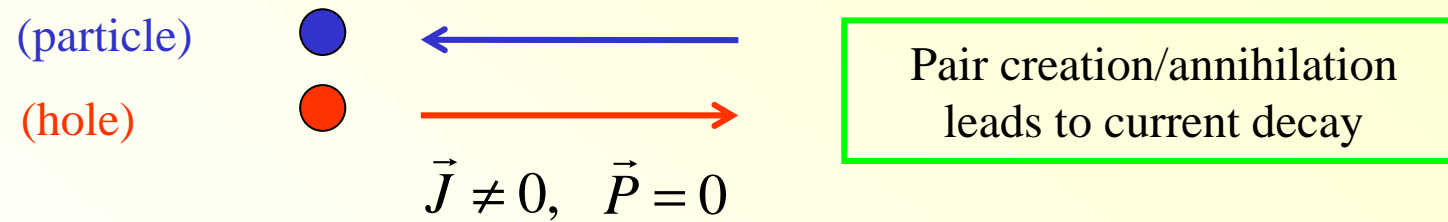
Collision-limited conductivity

Damle, Sachdev, (1996).

Fritz et al. (2008), Kashuba (2008)

Finite conductivity in a pure system at particle-hole symmetry ($\rho = 0$)!

- Key: Charge current without momentum



- **Finite collision-limited conductivity!** $\sigma(\mu = 0) < \infty$; $\sigma(\mu \neq 0) = \infty$
- Only marginal irrelevance of Coulomb:
Maximal possible relaxation rate,
set only by temperature

$$\tau_{ee}^{-1} \approx \alpha^2 \frac{k_B T}{\hbar}$$

→ Nearly universal conductivity at strong coupling

$$\sigma_{Drude} = \frac{e}{m} \rho \tau \rightarrow \sigma(\mu = 0) \sim \frac{e}{k_B T / v^2} \left(e \frac{(k_B T)^2}{(\hbar v)^2} \right) \frac{\hbar}{\alpha^2 k_B T} \sim \frac{1}{\alpha^2} \frac{e^2}{h}$$

Expect saturation as $\alpha \rightarrow 1$, and eventually phase transition to insulator

Marginal irrelevance of Coulomb:

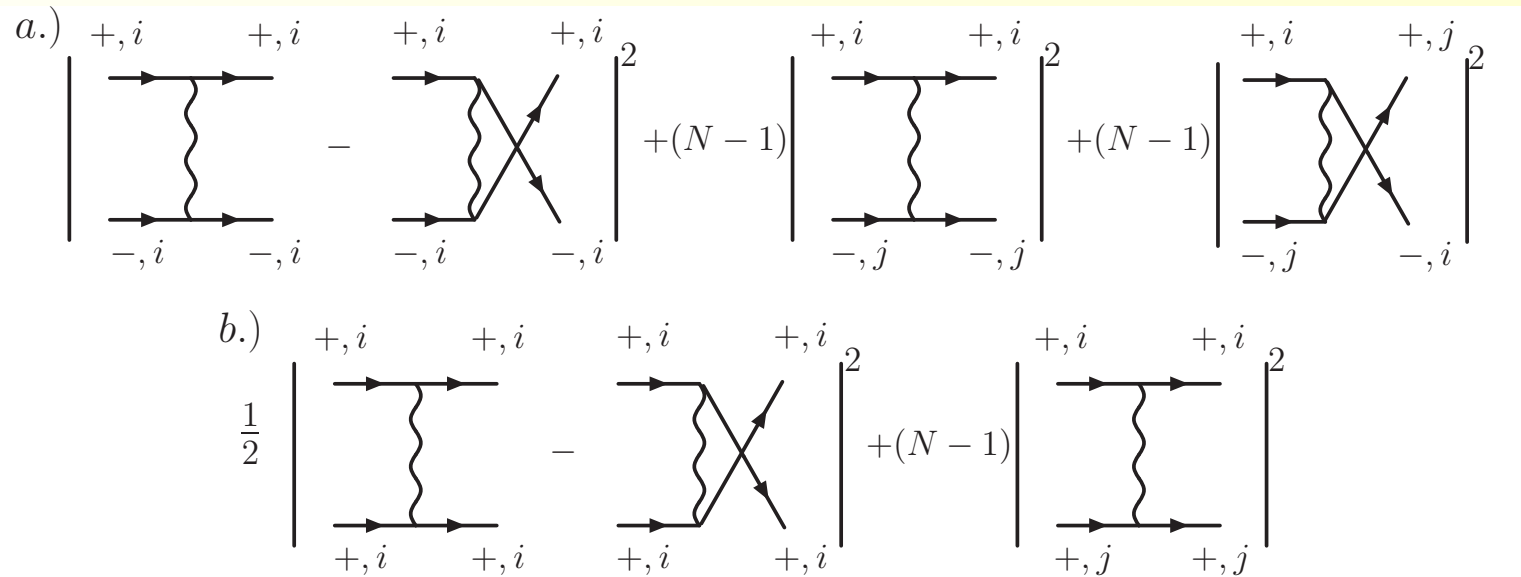
$$\alpha \approx \frac{4}{\log(\Lambda/T)} < 1$$

Boltzmann approach

L. Fritz, J. Schmalian, MM, and S. Sachdev, PRB 2008

Boltzmann equation in Born approximation

$$\left(\partial_t + e\mathbf{E} \cdot \frac{\partial}{\partial \mathbf{k}} \right) f_{\pm}(\mathbf{k}, t) = I_{\text{coll}}^{Cb}[\mathbf{k}, t | \{f_{\pm}(\mathbf{k}', t)\}] \propto \alpha^2(T)$$



→ Collision-limited conductivity:

$$\sigma(\mu = 0) = \frac{0.76 e^2}{\alpha^2(T) h}$$

Beyond weak coupling
approximation:

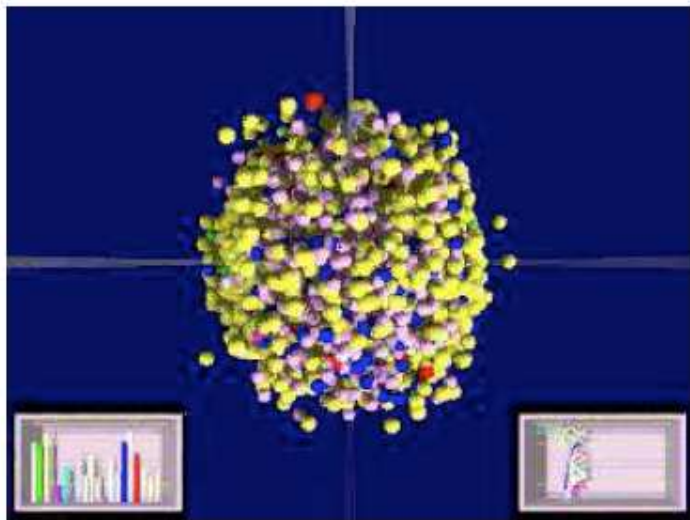
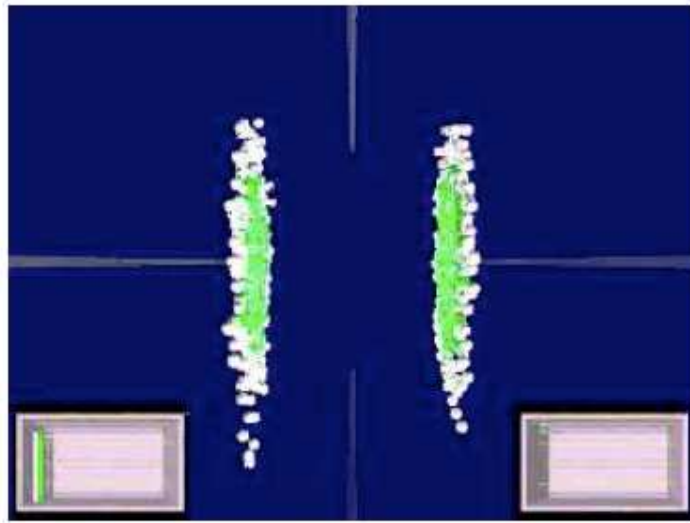
Graphene



Very strongly coupled, critical
relativistic liquids?

AdS – CFT !

Au+Au collisions at RHIC



Quark-gluon plasma is described
by QCD (nearly conformal,
critical theory)

—

Low viscosity fluid!

Compare graphene to Strongly coupled relativistic liquids

S. Hartnoll, P. Kovtun, MM, S. Sachdev (2007)

Obtain exact results via string theoretical AdS–CFT correspondence

→ Response functions for particular strongly coupled relativistic fluids
(for maximally supersymmetric SU(N) Yang Mills theory with $N \rightarrow \infty$ colors)
By mapping to weakly coupled gravity problem:

AdS - CFT [SU(N \gg 1)]

weak coupling - strong coupling

SU(N) transport from AdS/CFT

Gravitational dual to SUSY SU(N)-CFT₂₊₁: Einstein-Maxwell theory

$$I = \frac{1}{g^2} \int d^4x \sqrt{-g} \left[-\frac{1}{4} R + \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{3}{2} \right].$$

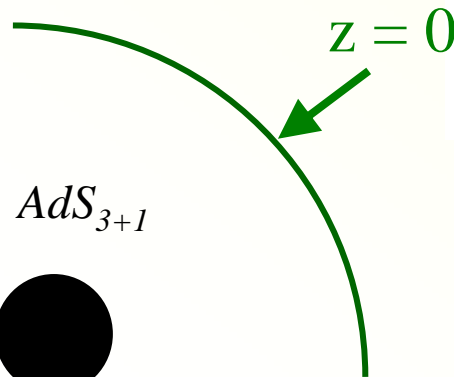
(embedded in M theory as $AdS_4 \times S^7$: $1/g^2 \sim N^{3/2}$)

It has a black hole solution (with electric and magnetic charge):

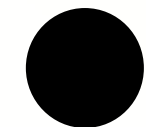
$$ds^2 = \frac{\alpha^2}{z^2} [-f(z)dt^2 + dx^2 + dy^2] + \frac{1}{z^2} \frac{dz^2}{f(z)},$$

$$F = h\alpha^2 dx \wedge dy + q\alpha dz \wedge dt,$$

$$f(z) = 1 + (h^2 + q^2)z^4 - (1 + h^2 + q^2)z^3.$$



AdS_{3+1}



Black hole

Electric charge q and magnetic charge, h

$\leftrightarrow \mu$ and B for the CFT

SU(N) transport from AdS/CFT

Gravitational dual to SUSY SU(N)-CFT₂₊₁: Einstein-Maxwell theory

$$I = \frac{1}{g^2} \int d^4x \sqrt{-g} \left[-\frac{1}{4} R + \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{3}{2} \right].$$

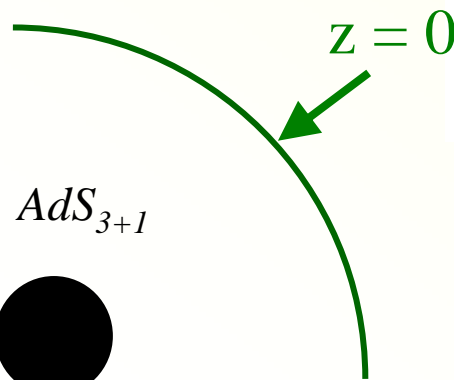
(embedded in M theory as $AdS_4 \times S^7$: $1/g^2 \sim N^{3/2}$)

It has a black hole solution (with electric and magnetic charge):

$$ds^2 = \frac{\alpha^2}{z^2} \left[-f(z) dt^2 + dx^2 + dy^2 \right] + \frac{1}{z^2} \frac{dz^2}{f(z)},$$

$$F = h\alpha^2 dx \wedge dy + q\alpha dz \wedge dt,$$

$$f(z) = 1 + (h^2 + q^2)z^4 - (1 + h^2 + q^2)z^3.$$



Black hole

Background \leftrightarrow Equilibrium

Transport \leftrightarrow Perturbations in $g_{tx,ty}, A_{x,y}$.

Response via Kubo formula from $\delta^2 I / \delta(g, A)^2$.

Compare graphene to Strongly coupled relativistic liquids

S. Hartnoll, P. Kovtun, MM, S. Sachdev (2007)

Obtain exact results via string theoretical AdS–CFT correspondence



- Confirm the structure of the hydrodynamic response functions $\sigma(\omega)$ etc.
- Allow to calculate the transport coefficients for a strongly coupled theory!

$$\text{SUSY - SU(N): } \sigma(\mu = 0) = \sqrt{\frac{2}{9}} N^{3/2} \frac{e^2}{h}$$

Compare graphene to Strongly coupled relativistic liquids

S. Hartnoll, P. Kovtun, MM, S. Sachdev (2007)

Obtain exact results via string theoretical AdS–CFT correspondence



- Confirm the structure of the hydrodynamic response functions $\sigma(\omega)$ etc.
- Allow to calculate the transport coefficients for a strongly coupled theory!

$$\text{SUSY - SU(N): } \sigma(\mu = 0) = \sqrt{\frac{2}{9}} N^{3/2} \frac{e^2}{h} \quad ; \quad \frac{\eta_{shear}}{s}(\mu = 0) = \frac{1}{4\pi} \frac{\hbar}{k_B}$$

Graphene – a nearly perfect liquid!

MM, J. Schmalian, and L. Fritz, (PRL 2009)



Anomalously low viscosity (like quark-gluon plasma)

“Heisenberg”

$$\frac{\eta}{s} \sim E_{qp} \tau \geq 1$$

→ Measure of strong coupling:

Conjecture from
AdS-CFT:

$$\frac{\text{shear viscosity}}{\text{entropy density}} = \frac{\eta}{s} > \frac{\hbar}{k_B} \frac{1}{4\pi}$$

Graphene – a nearly perfect liquid!

MM, J. Schmalian, and L. Fritz, (PRL 2009)



Anomalously low viscosity (like quark-gluon plasma)

“Heisenberg”

$$\frac{\eta}{s} \sim E_{qp} \tau \geq 1$$

→ Measure of strong coupling:

Conjecture from
AdS-CFT:

$$\frac{\text{shear viscosity}}{\text{entropy density}} = \frac{\eta}{s} > \frac{\hbar}{k_B} \frac{1}{4\pi}$$

Doped Graphene &
Fermi liquids:
(Khalatnikov etc)

$$\eta \propto n \cdot m v^2 \cdot \tau \rightarrow n \cdot E_F \cdot \frac{\hbar E_F}{(k_B T)^2}$$

$$s \propto k_B n \frac{T}{E_F}$$

$$\frac{\eta}{s} \sim \frac{\hbar}{k_B} \left(\frac{E_F}{T} \right)^3$$

Graphene – a nearly perfect liquid!

MM, J. Schmalian, and L. Fritz, (PRL 2009)



Anomalously low viscosity (like quark-gluon plasma)

“Heisenberg”

$$\frac{\eta}{s} \sim E_{qp} \tau \geq 1$$

→ Measure of strong coupling:

Conjecture from
AdS-CFT:

$$\frac{\text{shear viscosity}}{\text{entropy density}} = \frac{\eta}{s} > \frac{\hbar}{k_B} \frac{1}{4\pi}$$

Doped Graphene &
Fermi liquids:
(Khalatnikov etc)

$$\eta \propto n \cdot m v^2 \cdot \tau \rightarrow n \cdot E_F \cdot \frac{\hbar E_F}{(k_B T)^2}$$

$$s \propto k_B n \frac{T}{E_F}$$

$$\frac{\eta}{s} \sim \frac{\hbar}{k_B} \left(\frac{E_F}{T} \right)^3$$

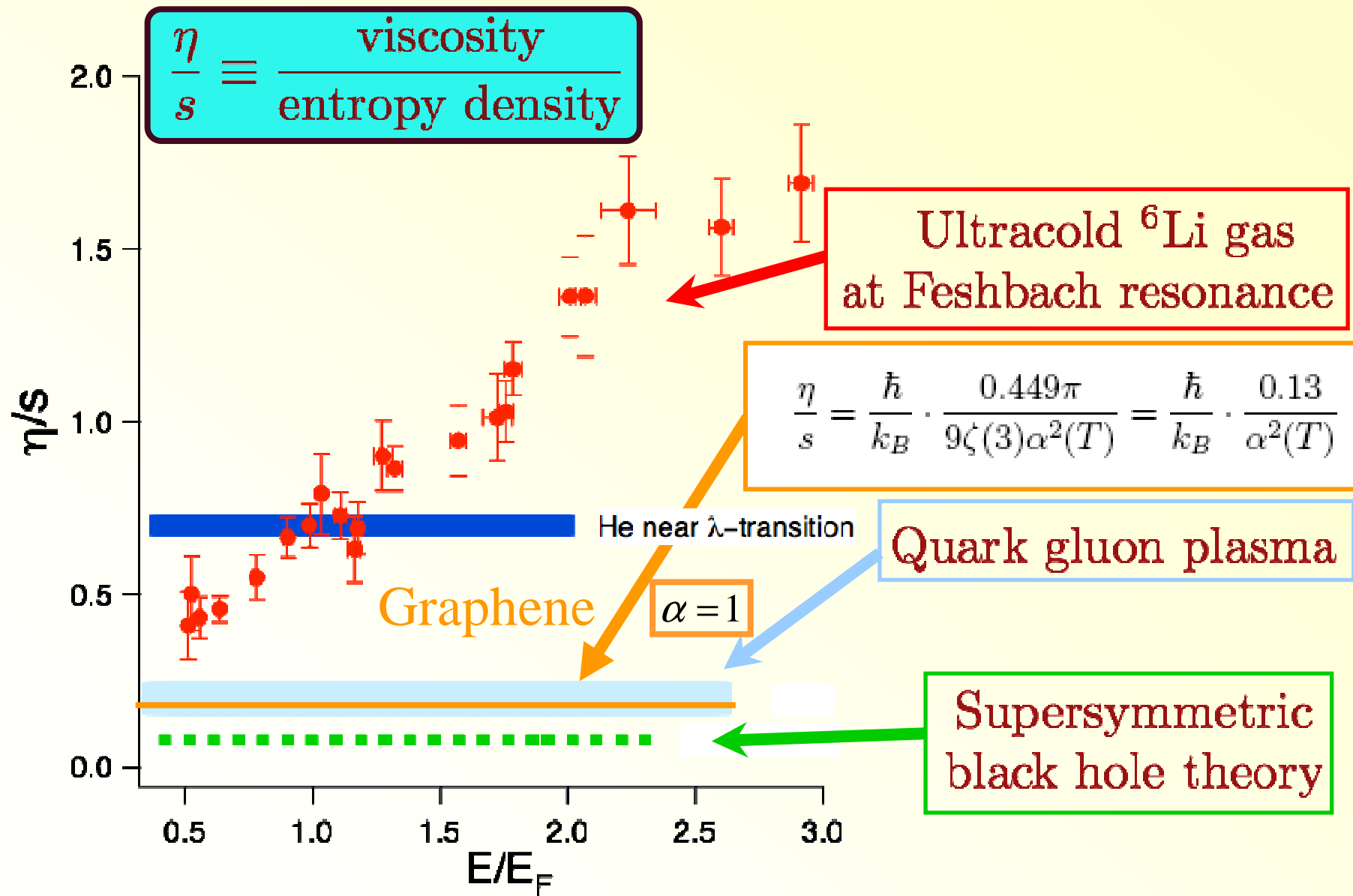
Undoped Graphene:

$$\eta \propto n \cdot m v^2 \cdot \tau \rightarrow n_{th} \cdot k_B T \cdot \frac{\hbar}{\alpha^2 k_B T} = \frac{\hbar}{\alpha^2} n_{th}$$

$$s \propto k_B n_{th}$$

Exact (Boltzmann-Born Approx):

$$\frac{\eta}{s} = \frac{\hbar}{k_B} \cdot \frac{0.449\pi}{9\zeta(3)\alpha^2(T)} = \frac{\hbar}{k_B} \cdot \frac{0.13}{\alpha^2(T)}$$



T. Schäfer, Phys. Rev. A 76, 063618 (2007).

A. Turlapov, J. Kinast, B. Clancy, Le Luo, J. Joseph, J. E. Thomas, J. Low Temp. Phys. 150, 567 (2008)

Electronic consequences of low viscosity?

MM, J. Schmalian, L. Fritz, (PRL 2009)

Electronic turbulence in clean, strongly coupled graphene?
(or at quantum criticality!)

Reynolds number:

$$\text{Re} = \frac{s/k_B}{\eta/\hbar} \times \frac{k_B T}{\hbar v/L} \times \frac{u_{\text{typ}}}{v}$$

Electronic consequences of low viscosity?

MM, J. Schmalian, L. Fritz, (PRL 2009)

Electronic turbulence in clean, strongly coupled graphene?
(or at quantum criticality!)

Reynolds number:

$$\text{Re} = \frac{s/k_B}{\eta/\hbar} \times \frac{k_B T}{\hbar v/L} \times \frac{u_{\text{typ}}}{v}$$

Strongly driven mesoscopic systems: (Kim's group [Columbia])

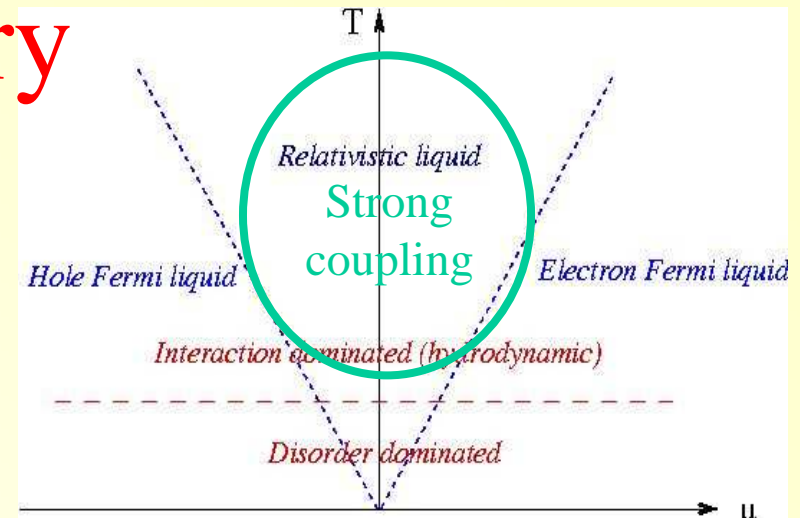
$$\begin{aligned} L &= 1\mu\text{m} \\ u_{\text{typ}} &= 0.1\text{v} \\ T &= 100\text{K} \end{aligned}$$

 $\text{Re} \sim 10 - 100$

Complex fluid dynamics!
(pre-turbulent flow)

New phenomenon in an
electronic system!

Summary



- Undoped graphene is strongly coupled in a large temperature window!
- Nearly universal strong coupling features in transport; many similarities with strongly coupled critical fluids (described by AdSCFT)
- Emergent relativistic hydrodynamics at low frequency
- Graphene: Nearly perfect quantum liquid!
→ Possibility of complex (turbulent?) current flow at high bias